

FORCE-FREE MAGNETOSPHERE OF AN ALIGNED ROTATOR



ANDREY TIMOKHIN (*Sternberg Astronomical Institute, Moscow, Russia*)

Pulsar Equation

Force-free magnetosphere of an aligned rotator rotating with the angular velocity Ω can be well described by the so-called pulsar equation

$$(\beta^2 x^2 - 1)(\partial_{xx}\psi + \partial_{zz}\psi) + \frac{\beta^2 x^2 + 1}{x} \partial_x \psi - S \frac{dS}{d\psi} + x^2 \beta \frac{d\beta}{d\psi} (\nabla\psi)^2 = 0. \quad (1)$$

Here all coordinates are normalized to $R_{LC}^{\text{cor}} \equiv c/\Omega$, corresponding to the Light Cylinder radius for co-rotating plasma. $\beta \equiv \Omega_F/\Omega$ is the normalized angular velocity of plasma rotation. We used such normalization for ψ that the dipole magnetic field function near the NS surface is given by

$$\psi^{\text{dip}} = \frac{x^2}{(x^2 + z^2)^{3/2}}, \quad (2)$$

The normalized poloidal current function $S \equiv (4\pi/c)(R_{LC}^2/\mu) I$, $\mu \equiv B_0 R_{NS}^3/2$. Magnetic field is expressed through the functions $\Psi \equiv \mu/R_{LC} \psi$ and I :

$$\mathbf{B} = \frac{\nabla\Psi \times \mathbf{e}_\phi}{\varpi} + \frac{4\pi I}{c \varpi} \mathbf{e}_\phi \quad (3)$$

At the Light Cylinder, $R_{LC} \equiv c/\Omega_F$, the pulsar equation has the form

$$2\beta \partial_x \psi = S \frac{dS}{d\psi} - \frac{1}{\beta} \frac{d\beta}{d\psi} (\nabla\psi)^2. \quad (4)$$

Any smooth solution must satisfy this equation at the Light Cylinder (LC). Angular velocity of plasma rotation in the open field line domain is

$$\Omega_F = \Omega + c \frac{\partial V}{\partial \Psi}. \quad (5)$$

where V is the non-corotational electric potential, caused by presence of an accelerating electric field in the polar cap of pulsar. If from a particularly theory of polar cap cascade we know β as a function of ψ , $\beta = \beta(\psi)$, the position of the Light Cylinder $R_{LC}(x, z) = c/\Omega_F[\psi(x, z)]$ must be determined self-consistently together with the solution of the pulsar equation.

Solution method

We assume Y-configuration of the magnetosphere, i.e. the existence of an equatorial current sheet with the return current. Equation (1) is solved numerically in a domain $x_{NS} \leq x \leq x_{\text{max}}$, $z_{NS} \leq z \leq z_{\text{max}}$

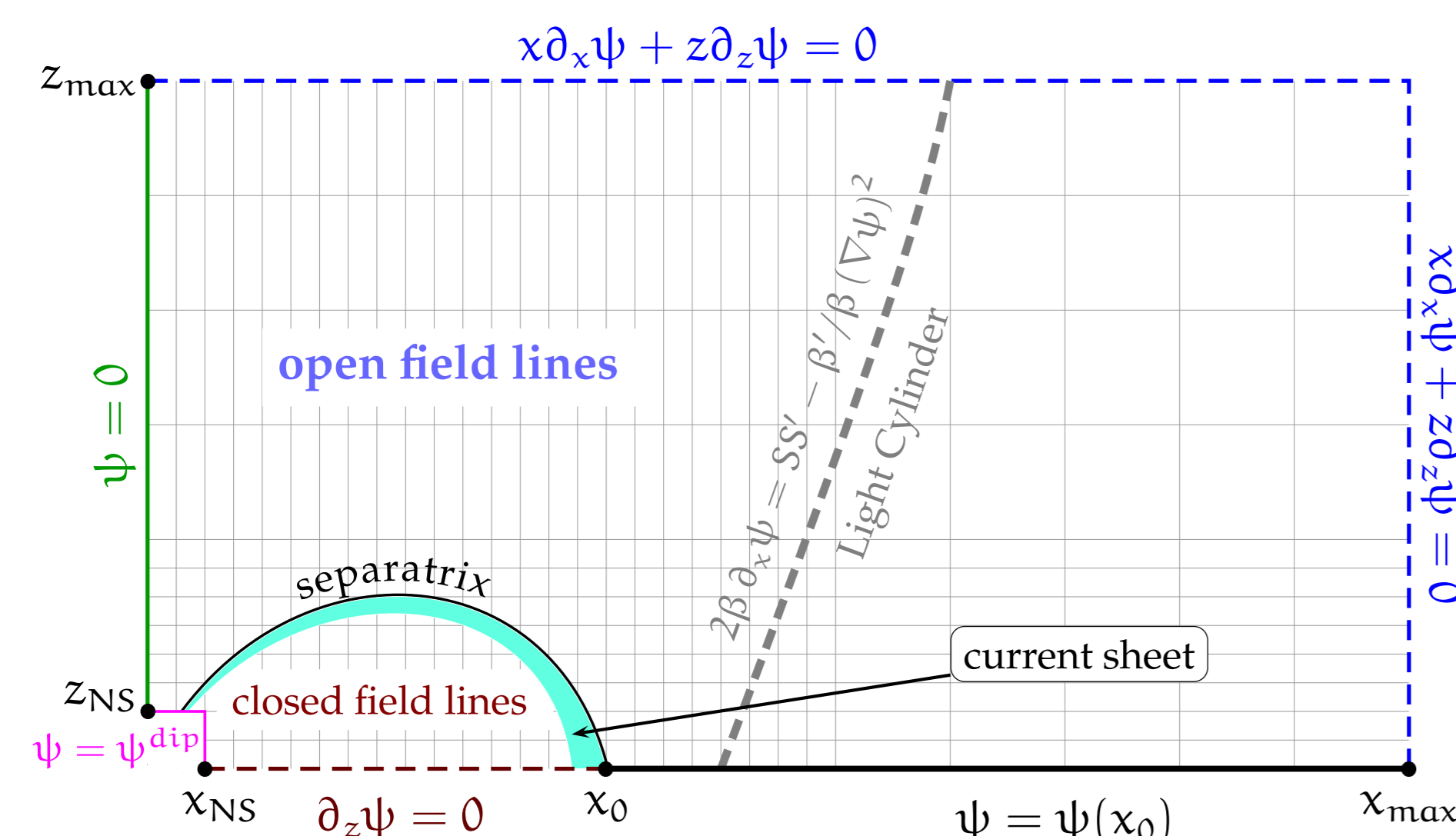


FIGURE 1: Calculation domain and boundary conditions

The return current flowing along the last closed magnetic field line is smeared over the region $[\psi_0 - d\psi, \psi_0]$.

Boundary conditions are shown in Fig. 1. The equation (1) is solved by a full multigrid (V-cycles) FAS scheme. As a smoother we used Gauss-Seidel scheme.

At each iteration step we find position of the LC $\{(x_{LCj}, z_j), j = 0 \dots n\}$ by solving numerically the equation

$$x_{LCj} = 1/\beta[\psi(x, z_j)]. \quad (6)$$

by the Newton method for each z -axis grid point z_j , and find poloidal current function $SS' \equiv S(dS/d\psi)$ from the equation (4). Then we use piece-polynomial interpolation for SS' and calculate $SS'(x, z) = SS'[\psi(x, z)]$ in each domain point.

β in the current sheet smoothly changes to the value $\beta = 1$ in the closed field line zone.

References

1. A. N. Timokhin, 2006, MNRAS (in press, doi: 10.1111/j.1365-2966.2006.10192.x), astro-ph/0511817
2. A. N. Timokhin, 2006, in preparation

Main Results

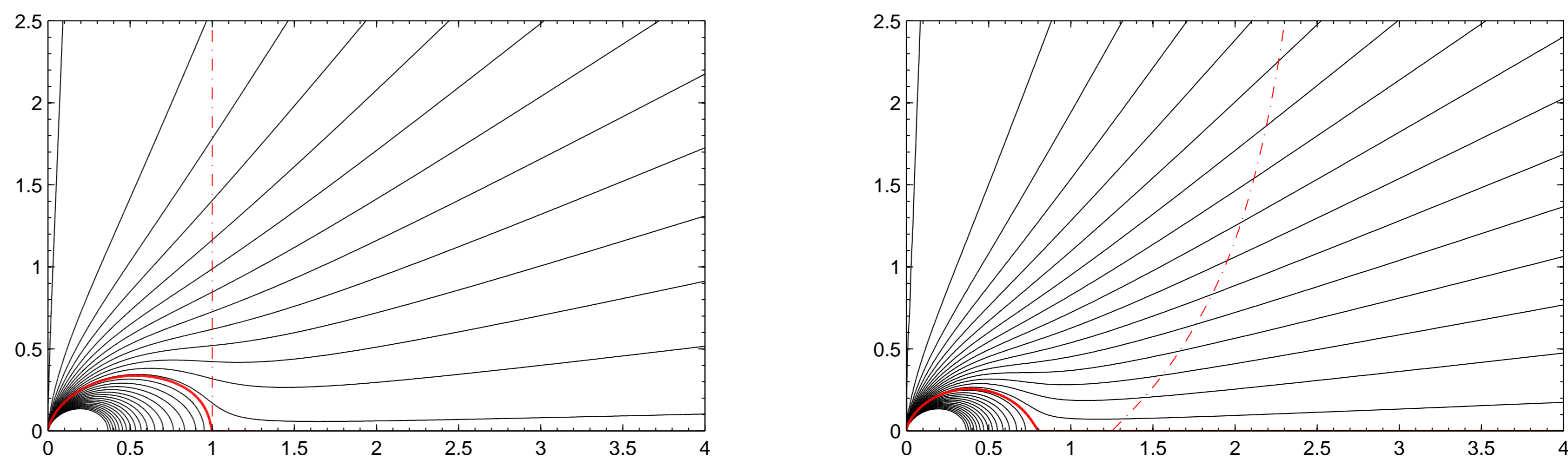


FIGURE 2: Structure of the pulsar magnetosphere: **left** – for $\beta \equiv 1$ and $x_0 = 1$; **right** – for $x_0 = 0.8$ and β varying in such a way that the current density in the polar cap of pulsar is nearly constant (see Fig.4 (right)). The Light Cylinder is shown by the dot-dashed line.

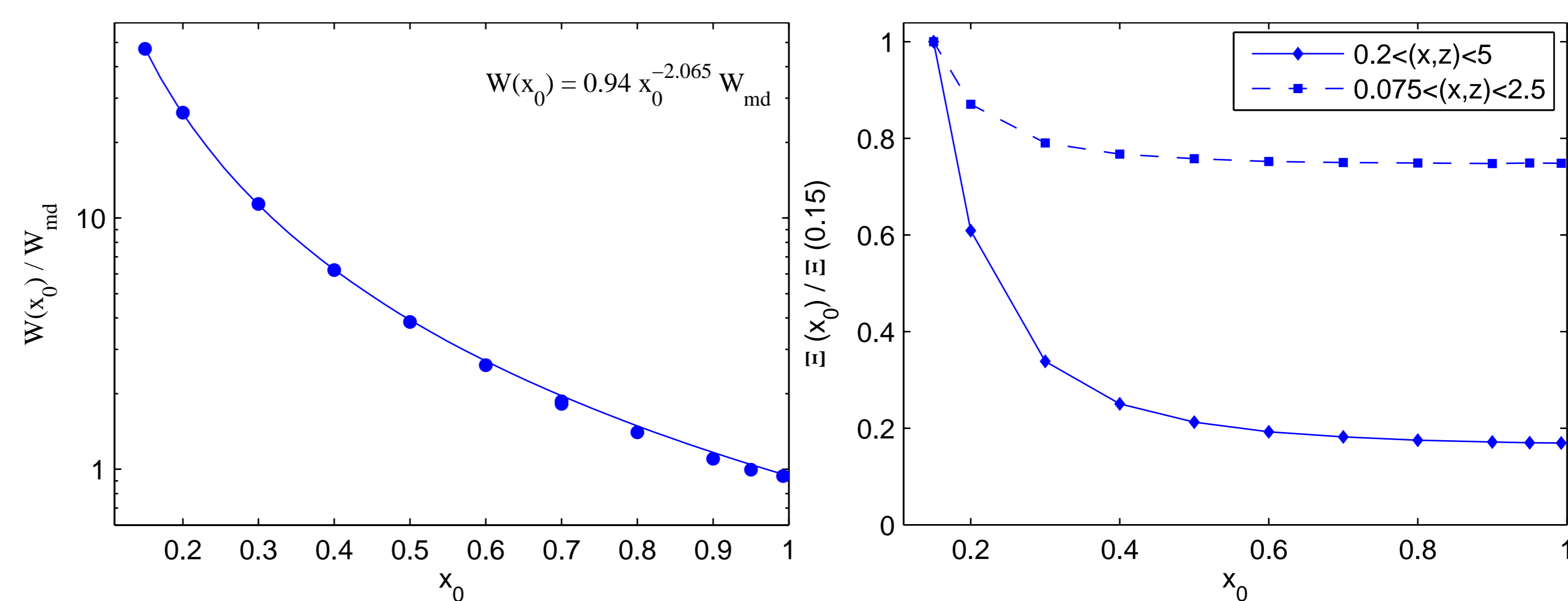


FIGURE 3: For $\beta \equiv 1$: **left** – energy losses normalized to the magnetodipolar energy losses as a function of x_0 ; **right** – total electromagnetic energy in two different volumes as a function of x_0

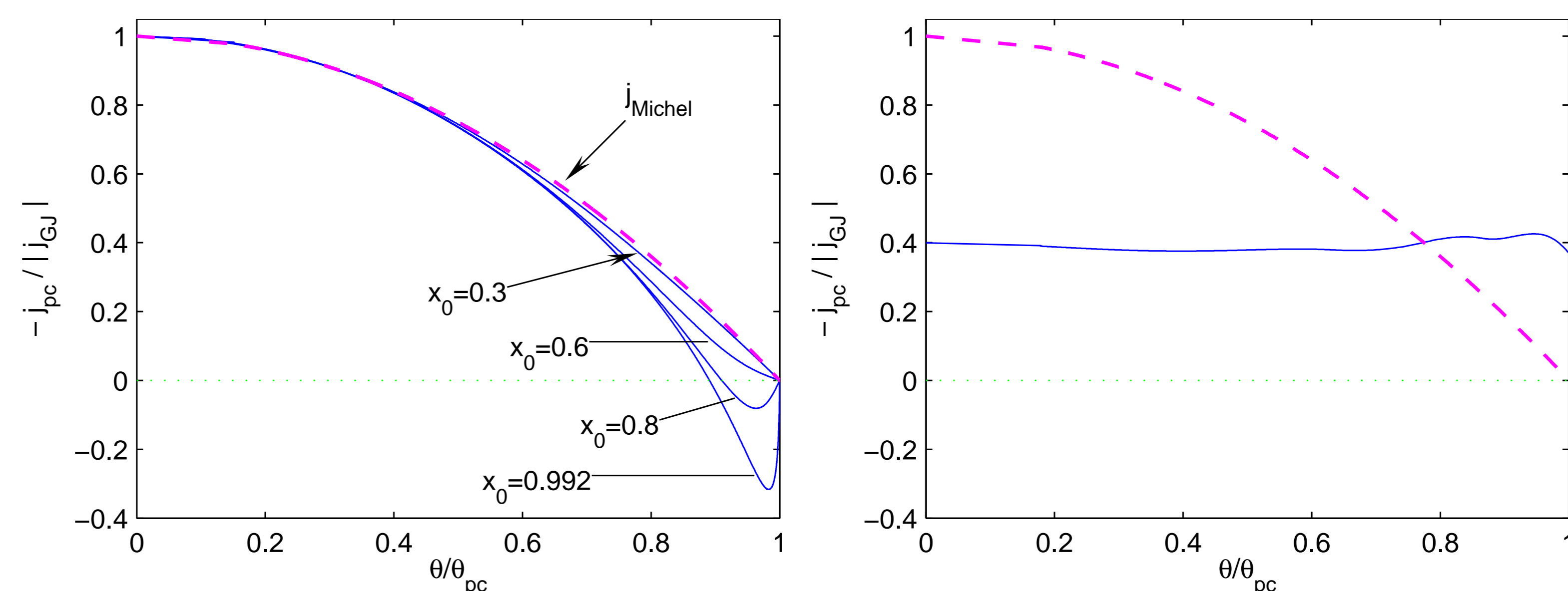


FIGURE 4: Poloidal current density in the polar cap of pulsar normalized to the Goldreich-Julian current density: **left** – for $\beta \equiv 1$ and different values of x_0 ; **right** – for variable β producing nearly constant current density. Michel current density is shown by the dashed line

Constant $\beta \equiv 1$

- Each solution has been checked for applicability of the force-free condition $E < B$. In none of them this condition is violated.
- The total energy of the electromagnetic field in the magnetosphere $\Xi = \int_{Vol} (B^2 + E^2)/(8\pi) dV$ decreases with increasing of x_0 , see Fig. 3 (right).
- Energy losses of the pulsar increase with decreasing of x_0 and are given by (see Fig. 3 (left))

$$W = 0.94 x_0^{-2.065} \times \frac{\mu^2 \Omega^4}{c^3}, \quad (7)$$

- In configurations with $x_0 > 0.6$ there is a volume return current flowing along the open magnetic field lines, which makes however only a small part of the whole return current. The current density (in units of j_{GJ}) close to the polar cap boundary increases with increasing of x_0 . The current density does not exceed the Goldreich-Julian current density $j_{GJ} \equiv \rho_{GJ} c$ and at most field lines it is less than j_{GJ}

The problem: current density distribution for $\beta \equiv 1$ is not compatible with the Space Charge Limited Flow regime of the stationary polar cap cascades. Current density distribution with return volume current is not compatible with any cascade model.

Variable β – magnetosphere with differential rotation in the open field line zone

- Current density can deviate from distributions shown in Fig. 4 (left) and could be made even nearly constant over the polar cap of pulsar. The latter would require adjusting of the plasma angular velocity, i.e. large values of the non-corotational electric potential V . Moreover, in this case the current density j is always less than j_{GJ} , see Fig. 4

Conclusions

- The magnetosphere of pulsar should evolve with time, i.e the relative size (in R_{LC}) of the closed field line zone should change. This will result in pulsar breaking index different from 3.
- The magnetosphere could rotate differentially – additional degree of freedom for adjusting of the current density required by the global structure of the magnetosphere to the current density produced by the polar cap cascades.
- Electromagnetic cascades in the polar cap of pulsar should be non-stationary.